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Cosmology and structure formation

Part III

Problems "solved" by Inflation

- ① flatness
- ② horizons
- ③ monopoles
- ④ Structure Problem

For a galaxy cluster · mass $10 M_\odot \sim 10^{15} M_p$

Quantum fluctuations in a non-inflating $\frac{\delta\phi}{c^2} \sim 10^{-5}$
Universe cannot give this amplitude $1/\sqrt{N} \sim 10^{-36}$

$$a \sim e^{Ht} \quad P = P_{\text{vac}} = \frac{3H^2}{8\pi G} = \text{const}$$

During inflation Universe is approx time-invariant

\Rightarrow quantum fluctuation amplitudes are time-INV.

\rightarrow fluctuations which "exit" the horizon at different times during inflation will have (statistically) the same amplitude but scales which differs by $\frac{\lambda_1}{\lambda_2} \sim \exp[H(t_1 - t_2)]$

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At much later time $\lambda \frac{a(t)}{a_0} = ct$ with the same amplitude the fluctuations reenter the horizon

If fluctuations originated as zero-point fluctuations of inflaton field $\rightarrow \frac{S\phi}{c^2}$ const on all scales as flutns cross horizon
Harrison-Zel'dovich fluctuations

⑨ Thermal history of Universe

T : t

10^{15} GeV 10^{-36} s Possible epoch of inflation (?) Generation of structure

$> 100 \text{ MeV}$ $< 10^{-4}$ s "quark soup" no bound hadrons. baryogenesis (?)

$\sim 100 \text{ MeV}$ 10^{-4} s quark-hadron phase transition

$\sim 3 \text{ MeV}$ 0.1 s ν 's "decouple" (ie rate for $e^- + e^+ \leftrightarrow \nu + \bar{\nu}$) becomes long compared to $t \sim 1/H$ Universe becomes transparent to ν 's

T t

$\sim 1 \text{ MeV}$ 1 s rates $P + e^- \xleftrightarrow{\text{co}} n + \bar{\nu}_e$ $n + e^+ \xleftrightarrow{\text{co}} p + \bar{\nu}_e$
for drop below H \rightarrow final $\frac{n}{P} \sim \exp\left(-\frac{\Delta m}{kT}\right) \sim 20\%$

0.3 MeV 10 s $e^+ e^-$ annihilation \rightarrow heats photon gas establishes
final $\frac{n_r}{n_s} \sim \frac{11}{3}$ $\frac{n_x}{n_{\text{baryon}}} \sim 10^9$

0.1 MeV 100 s $p + n \rightarrow D$ becomes possible Almost all n are
bound to ${}^4\text{He}$

0.05 MeV 400 s Nucleosynthesis complete 76% H, 24% ${}^4\text{He}$ + small
amounts of ${}^2\text{H}$, ${}^3\text{He}$, ${}^7\text{Li}$
Exact amounts depend strongly on $\frac{n_{\text{baryon}}}{n_r}$

⑨ Thermal history of Universe

T t

3×10^6 K 1 yr

Inelastic scattering of γ becomes inefficient
Black-body spectrum cannot be created after this time
Observed CMB \rightarrow no sign. input of γ after this

20,000 K 8000 yr

Universe switches from radiation $\xrightarrow{t \sim 10^4}$ matter domination.

3000 K 3×10^5 yr.

Recombination rate $e^+e^- \rightarrow H$ dominates reionization
by the Wien tail of the CMB photons. Recombination

273 K 132 Gyr TODAY

Nonlinear growth of structure in the standard Λ CDM paradigm with inflationary fluctuations adjusted to fit the WMAP data on CMB fluctuations

Evolution of the dark matter distribution in a thin slice centred on a galaxy cluster at $z=0$

Zoom into a rich cluster from a very large-scale image of the $z=0$ DM distribution in a thin slice

Evolution of the DM and gas distributions, the X-ray and S-Z images of a rich galaxy cluster

⑩ Structure formation

$$\rho(\varepsilon, t) = \bar{\rho}(t)(1 + S(\varepsilon, t)) \quad S \ll 1$$

Linearise the evolution eqns to get a linear equation
for $S(\varepsilon, t)$

$$FT, \text{ the spatial dependence} \quad S_k \sim \int d^3x S(\varepsilon, t) e^{ikx}$$

To get ODE's in time for $S_k(t)$

Result At late times and/or large scales in k $S_k(t) = b(t) S_k(t_0)$ b indep. of k $b \propto t^{2/3}$

⑨ Thermal history of Universe

Characterise linear fluctuations by their power spectrum $P(k)$

Inflationary fluct. are completely characterised by $P(k)$

$$\rightarrow P(k) \propto k^n \text{ with } n \approx 1$$

Waves with diff k have uncorrelated phases $\rightarrow \delta$ is a Gaussian random field

Shape of $P(k)$ at recombination depends

① Structure generator

② Nature of DM